Multiframe isochloric solution to υ-wire Transmission Field equations under Austwurnian Conditions (and implications thereof)

M. Hearn†
Department of Theoretical VX
VX Institute, Germany

***Abstract*—This paper presents a new solution to the υ-wire Transmission Field equations.**

We first introduce a new isochronous frame of reference to the nonlinear Austwurnian eigenproblem. We then prove that a natural abstract topological group that describes this system exists. Unlike previous attempts[1,3,6,7], this underlying theory is lossless and we also prove that a suitable inversion exists.

Utilising this theory we show several new results including a hypothesis that all VX-observable physical, imaginary and virtual fields arise from a single tristate field in ten-dimensional topological space. We also show that this space must have at most two timelike dimensions by eliminating the possibility of three or more timelike dimensions under Yalgeth's lemma. Further work is required to rule out the possibility of there being two timelike dimensions.

***Index Terms*—non-experimental, isochlorism, austwurnian conditions, hermian hypothesis, string theory, υ-wire fields**

I. Introduction

EXPERIMENTS suggest that what is typically observed at 0.88Δ is likely to be the self-interaction of the matter field ψ due to Connery's distortion. The question of whether it is composite or not has not been completely solved[2]; however there are quite strong clues[3,6,7] in favour of no additional internal field structure in the ninth vector as the $\psi\overline{\psi}$ coupling is typically nonzero. Under Austwurnian conditions however, the expected velocity of the local field $\langle v \rangle$ introduces an additional term $\delta_9(\phi, \frac{\partial\Delta}{\partial t})$. This sudden requirement to introduce of proper time t to the time-invariant field equations defies current understanding and complicates analysis. Previous work by Dalstaff *et al.* shows that this can lead to poles in the derivate $\psi\overline{\psi}$ field and suggests the possibility of nonlocal faster-than-light interactions (NFIs).

Experimental verification of this theoretical result has failed to be conclusive. Many think that it is unlikely for NFIs to exist and have continued theoretical research under this assumption. This has often lead to absurd conclusions that have to be corrected via renormalisation techniques borrowed from QM. This is an inelegant method. To quote Ryder[4], "despite the comparative success of renormalisation theory the feeling remains that there ought to be a more satisfactory way of doing things."

†M. Hearn is an independent student research associate of the VX Institute of Germany.

Draft generated October 10, 2014.

We discard renormalisation. We instead introduce a new virtual imaginary time parameter τ and rewrite the main υ-wire Transmission Field equations in terms of τ.

Thus the resulting higher order field has been shown to be aperiodic and lacking in translational symmetry. The Fourier transform of simulated γ emissions shows sharp delta peaks arranged in fivefold symmetric pattern.

II. Isochloric Methods

The central equation of the Harkonen subflux algorithm[5] is

$$\delta(g(x)) = \Sigma_i \frac{\delta(x - x_i)}{g'(x_i)}$$

This is an empirical equation; the measured agreement with isochloric wave function collapse is on the order of 3ns.

Instead of taking this equation as is, we shall now rederive this equation from scratch under modern VX theory. We shall first examine the zeroth-order deformation equation under a single frame:

$$(1 - q)\mathcal{L}[U(x; q) - u_0(x; q)] = c_0\, q\, \mathcal{N}[U(x; q)]$$

where the hyperparameter q varies as $0 \leq q \leq 1$. We shall take the convergence control parameter c_0 to satisfy the constraint

$$c_0 = g(c_0) \left| \Delta^{\phi} \right\rangle$$

This has traditionally been solved numerically. We shall now present an analytic solution to this equation by first demonstrating that it is linearisable. This is a lossy transformation (how to perfectly correct for this error is shown in section 5, *Renormalisation without Renormalisation*).

Under perturbation theory[12] one may assume c_0 is a perturbation from a set-point c_{0_0}. This approximation may be written in modern notation as

$$c_0(\epsilon) = c_{0_0} + \sum_i^{\infty} \alpha_i \epsilon^i$$

where ϵ is the perturbation parameter.

It is possible to further perturb this solution to produce $c_{0_{0_0}}$, $c_{0_{0_{0_0}}}$, *et cetera*. We can take the limit of an infinite number of perturbations and associate the stack of ϵs with a new parameter τ. The seventh Austwurnian constraint reduces the possible forms for ϵ to three forms that may be Morrington